



Program: RFEM 5, RSTAB 8

Category: Geometrically Linear Analysis, Isotropic Linear Elasticity, Elastic Foundation, Member

Verification Example: 0002 – Cantilever Beam on an Elastic Winkler Foundation

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Description

A cantilever beam of length L with rectangular cross-section of height h and width b lying on an elastic Winkler foundation of stiffness $C_{1,z}$ is loaded by a distributed loading q_z . Neglecting self-weight, determine the maximum deflection u_z and maximum bending moment M_y of the beam. Calculate the same example also for a plate of the same height and width as the cantilever.

Material	Isotropic Linear Elastic	Modulus of Elasticity	E	210.000	GPa
		Shear Modulus	G	105.000	GPa
Geometry	Cantilever	Length	L	4.000	m
		Height	h	0.200	m
		Width	b	0.005	m
Member Foundation	Winkler Elastic	Stiffness	$C_{1,z}$	500.000	kN/m ²
Plate Foundation		$C_{u,z} = \frac{C_{1,z}}{b}$	100000.000	kN/m ³	
Load	Member	Distributed	q_z	1.000	kN/m
	Plate	Distributed	$q = \frac{q_z}{b}$	200.000	kN/m ²

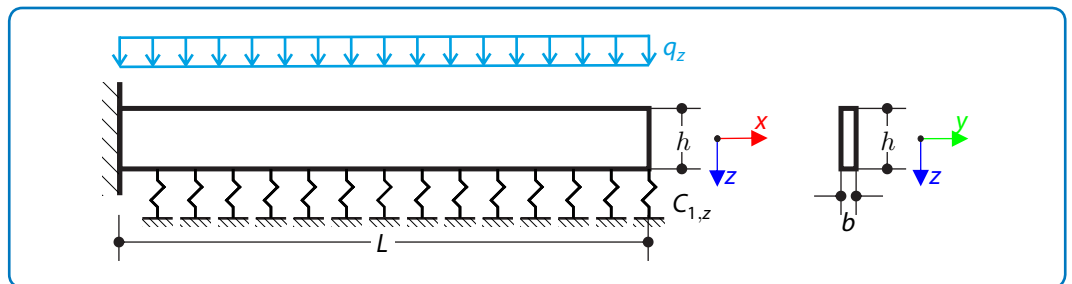


Figure 1: Problem sketch

Analytical Solution

Member Calculation

The governing differential equation for a beam on an elastic foundation can be expressed as

$$E I_y \frac{d^4 u_z}{dx^4} + C_{1,z} u_z = q_z \quad (2 - 1)$$

with the moment of inertia $I_y = \frac{1}{12} b h^3 = 3.33 \times 10^{-6} \text{ m}^4$. Dividing by $E I_y$ and setting $\beta^4 = \frac{C_{1,z}}{4 E I_y}$, equation (2 - 1) can be rewritten as

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$$\frac{d^4 u_z}{dx^4} + 4\beta^4 u_z = \frac{q_z}{EI_y} \quad (2-2)$$

The solution of (2-2) can be obtained as the superposition of the solutions of a particular integral, which can be expressed, assuming $u_z = C = \text{const}$, as

$$0 + 4\beta^4 C = \frac{q_z}{EI_y} = \text{const} \quad (2-3)$$

which leads to

$$C = \frac{q_z}{4\beta^4 EI_y} = \frac{q_z}{4 \frac{C_{1,z}}{4EI_y} EI_y} = \frac{q_z}{C_{1,z}} \quad (2-4)$$

and the solution of the characteristic equation

$$\frac{d^4 u_z}{dx^4} + 4\beta^4 u_z = 0 \quad (2-5)$$

To solve the characteristic equation (2-5), assume that $u_z = Ae^{\lambda x}$, hence

$$\lambda^4 + 4\beta^4 = 0 \quad (2-6)$$

Then the solution for λ can be expressed as

$$\lambda^4 = -4\beta^4 \Rightarrow \lambda_{k+1} = \sqrt[4]{(4\beta^4)} \left[\cos\left(\frac{\pi + 2k\pi}{4}\right) + i \sin\left(\frac{\pi + 2k\pi}{4}\right) \right] \quad (2-7)$$

where $k = 0, 1, 2, 3$. Equation (2-7) can be rewritten for all four variants as

$$\lambda_1(k=0) = \beta\sqrt{2} \left[\cos\left(\frac{\pi}{4}\right) + i \sin\left(\frac{\pi}{4}\right) \right] = \beta(1+i) \quad (2-8)$$

$$\lambda_2(k=1) = \beta\sqrt{2} \left[\cos\left(\frac{3\pi}{4}\right) + i \sin\left(\frac{3\pi}{4}\right) \right] = \beta(-1+i) \quad (2-9)$$

$$\lambda_3(k=2) = \beta\sqrt{2} \left[\cos\left(\frac{5\pi}{4}\right) + i \sin\left(\frac{5\pi}{4}\right) \right] = \beta(-1-i) \quad (2-10)$$

$$\lambda_4(k=3) = \beta\sqrt{2} \left[\cos\left(\frac{7\pi}{4}\right) + i \sin\left(\frac{7\pi}{4}\right) \right] = \beta(1-i) \quad (2-11)$$

Therefore, the solution u_z of (2-5) takes the form

$$u_z = \sum_{i=1}^4 A_i e^{\lambda_i x} \quad (2-12)$$

Substituting equations (2-8)–(2-11), (2-12) can be rewritten as

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$$u_z = A_1 e^{\beta x(1+i)} + A_2 e^{\beta x(-1+i)} + A_3 e^{\beta x(-1-i)} + A_4 e^{\beta x(1-i)} = \quad (2-13)$$

$$e^{\beta x} (A_1 e^{\beta xi} + A_4 e^{-\beta xi}) + e^{-\beta x} (A_2 e^{\beta xi} + A_3 e^{-\beta xi}) \quad (2-14)$$

Incorporating $e^{\beta xi} = \cos(\beta x) + i \sin(\beta x)$ into (2-13) yields

$$u_z = e^{\beta x} (C_1 \cos(\beta x) + C_2 \sin(\beta x)) + e^{-\beta x} (C_3 \cos(\beta x) + C_4 \sin(\beta x)) = \quad (2-15)$$

$$\cos(\beta x) (C_1 e^{\beta x} + C_3 e^{-\beta x}) + \sin(\beta x) (C_2 e^{\beta x} + C_4 e^{-\beta x}) \quad (2-16)$$

which can be further simplified using a new set of unknowns and the definition of hyperbolic functions

$$u_z = \cos(\beta x) (D_1 \cosh(\beta x) + D_2 \sinh(\beta x)) + \sin(\beta x) (D_3 \cosh(\beta x) + D_4 \sinh(\beta x)) \quad (2-17)$$

The final solution of equation (2-2) is constructed by the superposition of the solutions (2-4) and (2-17)

$$u_z = \cos(\beta x) (D_1 \cosh(\beta x) + D_2 \sinh(\beta x)) + \sin(\beta x) (D_3 \cosh(\beta x) + D_4 \sinh(\beta x)) + \frac{q_z}{C_{1,z}} \quad (2-18)$$

To obtain values for constants D_1 - D_4 , four cantilever boundary conditions have to be applied

$$1) \quad u_z(0) = 0 \quad (2-19)$$

$$2) \quad \frac{du_z}{dx}(0) = 0 \quad (2-20)$$

$$3) \quad M_y(L) = E I_y \frac{d^2 u_z}{dx^2}(L) = 0 \Rightarrow \frac{d^2 u_z}{dx^2}(L) = 0 \quad (2-21)$$

$$4) \quad V_z(L) = E I_y \frac{d^3 u_z}{dx^3}(L) = 0 \Rightarrow \frac{d^3 u_z}{dx^3}(L) = 0 \quad (2-22)$$

which leads to

$$1) \quad u_z(0) = D_1 + \frac{q_z}{C_{1,z}} = 0 \Rightarrow D_1 = -\frac{q_z}{C_{1,z}} \quad (2-23)$$

$$2) \quad \frac{du_z}{dx}(0) = \beta (D_2 + D_3) = 0 \Rightarrow D_2 = -D_3 \quad (2-24)$$

$$3) \quad \frac{d^2 u_z}{dx^2}(L) = -2\beta^2 (D_1 s s_h + D_2 s c_h - D_3 c s_h - D_4 c c_h) = 0 \quad (2-25)$$

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$$4) \quad \frac{d^3 u_z}{dx^3}(L) = -2\beta^3(D_1 s c_h + D_2 s s_h + D_1 c s_h + D_2 c c_h - D_3 c c_h - D_4 c s_h + D_3 s s_h + D_4 s c_h) = 0 \quad (2-26)$$

where $s = \sin(\beta L)$, $c = \cos(\beta L)$, $s_h = \sinh(\beta L)$, and $c_h = \cosh(\beta L)$. Substituting (2-23) and (2-24) into (2-25) and (2-26), the following relations are obtained

$$3) \quad -\frac{q_z}{C_{1,z}}(s s_h) - D_3(s c_h + c s_h) - D_4(c c_h) = 0 \quad (2-27)$$

$$4) \quad -\frac{q_z}{C_{1,z}}(s c_h + c s_h) - D_3(s s_h + c c_h + c c_h - s s_h) - D_4(c s_h - s c_h) = 0 \quad (2-28)$$

Combining (2-27), (2-28) yields

$$\begin{bmatrix} s c_h + c s_h & c c_h \\ 2c c_h & c s_h - s c_h \end{bmatrix} \begin{bmatrix} D_3 \\ D_4 \end{bmatrix} = \begin{bmatrix} -\frac{q_z}{C_{1,z}} s s_h \\ -\frac{q_z}{C_{1,z}}(s c_h + c s_h) \end{bmatrix} \quad (2-29)$$

Solving (2-29) leads to the coefficients D_3 and D_4 in the form

$$D_3 = -\frac{q_z}{C_{1,z}} \left(\frac{c s + s_h c_h}{2 + c_h^2} \right) \quad (2-30)$$

$$D_4 = -\frac{q_z}{C_{1,z}} \left(\frac{c^2 - c_h^2}{c^2 + c_h^2} \right) \quad (2-31)$$

Finally, substituting equations (2-23), (2-24), (2-30), and (2-31) into (2-18) and setting $x = L$, the value for the maximum deflection u_z is obtained

$$u_{z,\max} = u_z(L) = c(D_1 c_h + D_2 s_h) + s(D_3 c_h + D_4 s_h) + \frac{q_z}{C_{1,z}} = 2.498 \text{ mm} \quad (2-32)$$

Similarly, setting $x = 0$ and substituting (2-25) and (2-31) into (2-21) gives the value for the maximum bending moment M_y

$$M_{y,\max} = M_y(0) = -EI_y \frac{d^2 u_z}{dx^2}(0) = -2EI_y \beta^2 D_4 = -1.146 \text{ kNm} \quad (2-33)$$

Plate Calculation

The cantilever is also calculated using plate elements of width b and height h on a Pasternak foundation. The example yields the same numerical results, so the theory is identical. The parameter $C_{2,z}$ describing the Pasternak foundation for plates that yields the same results is equal to $C_{u,z} = \frac{C_{1,z}}{b} = 100000 \text{ kN/m}^3$.

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Note that, in order to approximate the member solution exactly, the Poisson ratio is zero.

RFEM 5 and RSTAB 8 Settings

- Modeled in version RFEM 5.16.01 and RSTAB 8.16.01
- The element sizes are $l_{FE} = 0.400$ m (member) and $l_{FE} = 0.100$ m (plate)
- Geometrically linear analysis is considered
- Isotropic linear elastic material model is used
- The Kirchhoff plate theory is used
- Shear stiffness of members is deactivated

Results

Structure File	Entity	Program
0002.01	Member	RFEM 5
0002.02	Member	RSTAB 8
0002.03	Plate	RFEM 5

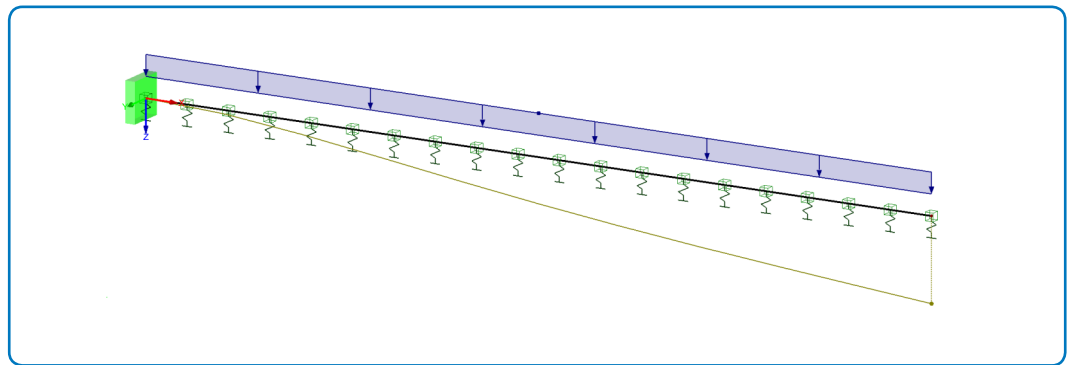


Figure 2: RFEM 5 Model – Member

As seen from the following comparisons, excellent agreement between the analytical solutions and numerical outputs has been achieved.

Analytical Solution	RFEM 5 (Member)		RSTAB 8 (Member)		RFEM 5 (Plate)	
$u_{z,max}$ [mm]	$u_{z,max}$ [mm]	Ratio [-]	$u_{z,max}$ [mm]	Ratio [-]	$u_{z,max}$ [mm]	Ratio [-]
2.498	2.498	1.000	2.498	1.000	2.495	0.999

Analytical Solution	RFEM 5 (Member)		RSTAB 8 (Member)		RFEM 5 (Plate)	
$M_{y,max}$ [kNm]	$M_{y,max}$ [kNm]	Ratio [-]	$M_{y,max}$ [kNm]	Ratio [-]	$m_{x,max} \times b$ [kNm]	Ratio [-]
-1.146	-1.146	1.000	-1.146	1.000	-1.139	0.994